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THE TEMPERATURE DEPENDENCE OF A LARGE DYNAMIC RANGE PHOTODETECTOR STRUCTURE (U)

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DEFENCE RESEARCH ESTABLISHMENT OTTAWA
TECHNICAL NOTE 91-35

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ABSTRACT

A recently developed photodetector circuit exploits the exponential voltage-to-current characteristic of a MOSFET operated in the subthreshold region to achieve a logarithmic steady state response. This paper analyzes the temperature dependence of the circuit operation and presents experimental results demonstrating the capabilities and limitations of the model.

RESUME

La réponse logarithmique en régime permanent d'un nouveau circuit pour photodétecteur est obtenue grâce à la relation exponentielle entre la tension et le courant d'un MOSFET opérant sous le seuil. Cette publication contient une analyse de l'opération du circuit en fonction de la température. Des résultats expérimentaux démontrent les capacités et les limitations du modèle présenté.

EXECUTIVE SUMMARY

A recently developed photodetector circuit exploits the exponential voltage-to-current characteristic of a MOSFET operated in the subthreshold region to achieve a logarithmic steady state response. This paper analyzes the temperature dependence of the circuit operation and presents experimental data which is generally consistent with the model. It is observed that the behaviour of the photodetector is similar in some respects to that of a photodiode operated in the solar cell mode, but has a greater sensitivity to fabrication process parameters. A test circuit which would facilitate accurate measurements at low illumination levels is proposed.

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1.0 INTRODUCTION

The dynamic range of photodetectors is important in many applications. One solution is to operate a photodiode in the solar cell mode. The steady state open circuit output voltage is given by

$$V = \frac{mkT}{q} \ln\left(\frac{I_p + I_o}{I_o}\right) \quad (1)$$

where m is a constant between 1 and 2, k is Boltzmann's constant, T is the absolute temperature in degrees Kelvin, q is the charge of an electron, I_o is the reverse saturation current of the photodiode and I_p is the photocurrent. Since the photocurrent is proportional to the intensity of the illumination, the steady state response is logarithmic for $I_p >> I_o$. Unfortunately, operation of photodiodes in the solar cell mode on a monolithic array requires dielectric isolation techniques to electrically isolate the photodiodes.

The recently developed photodetector circuit illustrated in Figure 1 uses the exponential voltage to current characteristic of a MOSFET operated in the subthreshold region to achieve a similar behaviour [1]. It is attractive for the implementation of monolithic arrays since it is compatible with standard MOS processing technology.

The temperature dependence of the circuit operation has not been considered in published analysis, but is potentially important in some applications. This is discussed in this report.

2.0 CIRCUIT MODEL

For unloaded steady state operation, the MOSFET source and photodiode currents, I_s and I_D respectively, satisfy the equality

$$I_S = I_D \quad (2)$$

With appropriate analytic behavioral models for I_D and I_s, this relationship can be evaluated to determine the circuit behaviour as a function of environmental and circuit parameters.

2.1 MOSFET Model

The behaviour of MOSFETs in the subthreshold region is of considerable significance in a number of applications. MOSFETs may be operated in the subthreshold region in low power analog circuits where the relatively high ratio of transconductance to

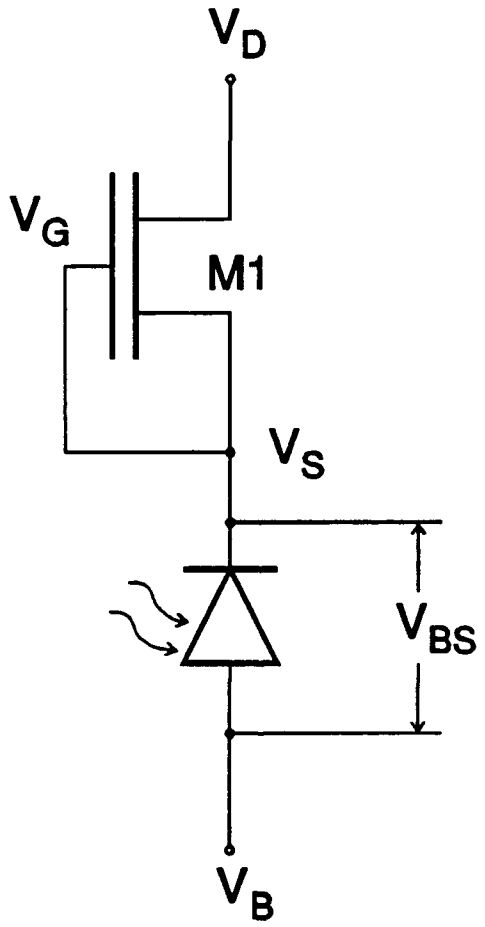

FIGURE 1: LARGE DYNAMIC RANGE PHOTODETECTOR CIRCUIT

drain current achievable is desirable [2]. Another application is found in analog neural networks which exploit the exponential relationship between gate-source bias and drain current [3]. However, most published analysis of the temperature dependence of MOSFETs operated in the subthreshold region concerns the design implications for digital VLSI circuits where it is important to ensure that a transistor biased in the OFF condition does not conduct a significant current over the operating temperature range [4].

The wide dynamic range photodetector circuit is analyzed using an expression for the current given in [5],

$$I_s = (W/L)\,DqN_DL_B(n_i/N_D)^2\exp(-\beta V_{BS})\exp(\beta\phi_{sat})\,[1-\exp(-\beta V_{DS})]\,(2\beta\phi_{sat})^{-\frac{1}{2}} \qquad (3)$$

where

L	=	channel length
W	=	channel width
D	=	diffusion constant $(kT/q)\mu$, where μ = carrier mobility
N_D	=	channel dopant impurity concentration
n_i	=	intrinsic carrier concentration α $T^{1.5}\exp(-E_g/2kT)$, where E_g = bandgap potential (~1.1 eV for Silicon)
β	=	$(\text{thermal voltage})^{-1} = q/kT$
V_{BS}	=	source to substrate voltage
V_{DS}	=	source to drain voltage
L_B	=	extrinsic Debye length = $[\epsilon_s\epsilon_o/\beta qN_D]^{1/2}$
ϵ_s	=	relative permittivity of Silicon = 11.7
ϵ_o	=	permittivity of free space
ϕ_{sat}	=	band bending potential at the channel - gate dielectric interface for channel pinch off.

Equation (3) can be solved explicitly as a function of the MOSFET terminal potentials since the term $\beta\phi_{sat}$ is a function of V_{BS} and of the effective source to gate voltage V_{GS}. The relationship is

$$\beta\phi_{sat} = \beta V_{GS} + \beta V_{BS}\,a^2/2 - a\,(\beta\,(V_{BS}+V_{GS}) + a^2/4)^{1/2} \qquad (4)$$

with the parameter a given by

$$a = \sqrt{2}\,(\epsilon_s/\epsilon_{ox})\,(t_{ox}/L_B)\ , \qquad (5)$$

where

t_{ox} = gate dielectric thickness
and ϵ_{ox} = relative permittivity of the gate dielectric (~ 3.7 for SiO_2).

Equation (3) can be rewritten to explicitly show the temperature dependence:

$$I_s = GT^{7/2} \frac{\left[\exp\left[(-E_g/kT) + (qV_{GS}/kT + \frac{a^2}{2}(T_o)\, T_o/T) \right.\right.}{\left.\left. -a(T_o)\, T_o/T \left(\frac{q}{kT_o}(V_{BS}+V_{GS}) + \frac{a^2(T_o)}{4} \right)^{\frac{1}{2}} \right] (1-\exp(-qV_{DS}/kT)) \right]} \qquad (6)$$
$$\sqrt{2} \left[\left(qV_{GS}/kT + qV_{BS}/kT + \frac{a^2}{2}(T_o)\,(T_o/T) \right) - a(T_o) \frac{T_o}{T} \left(q(V_{GS}+V_{BS})/kT_o + \frac{a^2}{4}(T_o) \right]^{\frac{1}{2}} \right)^{\frac{1}{2}} \right]$$

where

$$G = \frac{wD}{L} \left[\frac{k^5 \epsilon_s \epsilon_o}{N_D^3} \right]^{\frac{1}{2}}$$

contains parameters which are relatively insensitive to terminal voltages or temperature and $a(T_o)$ is the value of a at a temperature T_o. Although some parameters such as E_g and D are known to be dependent on temperature [6]-[7], the resultant error is not significant over the range of operating temperatures encountered in most practical applications.

For the wide dynamic range photodetector circuit, Equation (6) can be greatly simplified with minimal errors. First, the physical short circuit between the MOSFET gate and source terminals constrains V_{GS} to be equal to $-V_{FB}$. V_{FB} is the flatband voltage which is necessary to account for surface states, oxide charge and the metal-semiconductor potential ϕ_{MS}. Theoretical estimates of V_{FB} are very inaccurate, consequently, it is usually determined experimentally [5].

Second, the term $(1-\exp(-qV_{DS}/kT))$ in the numerator can be neglected if the circuit is biased for normal operation. Finally, as the gate to source threshold voltage V_{TS} of the MOSFET must be negative with $V_{BS} = 0$ for correct circuit operation, it can be shown that the upper limit for the parameter a is approximately 2.5 for the usual ranges of silicon process parameters [5]. Since (qV_{BS}/kT) will be much larger than a^2, the terms in Equation (6) containing a^2 can be neglected with minimal error.

The simplified equation for $I_S(T,V_{BS})$ is therefore given by

$$I_s(T,V_{BS}) = \frac{GT^{7/2}\exp[-E_g/kT+qV_{GS}/kT-a(T_o)(T_o/T)(q(V_{GS}+V_{BS})/kT_o)^{\frac{1}{2}}]}{\sqrt{2}(qV_{GS}/kT+qV_{BS}/kT-a(T_o)(T_o/T)(q(V_{GS}+V_{BS})/kT_o)^{\frac{1}{2}})^{\frac{1}{2}}} \quad (7)$$

Equation (3) for the MOSFET subthreshold current is derived from a one dimensional model of the MOSFET channel potential. The accuracy of this model is dependent on the channel length; as this is reduced, various additional phenomena degrade the validity of the model. These include the effects of the source and drain potentials on the MOSFET channel potential, transport effects such as carrier velocity saturation, and high electric field effects such as the injection of electrons into the gate oxide and thereby shifting the flatband voltage [5].

The effect of the MOSFET drain potential on the subthreshold current is useful in the photodetector circuit since it permits the DC operating point for V_{BS} to be modified by varying the drain voltage V_D. Furthermore, the dependence of the subthreshold current on the source potential V_{BS} provides additional compression which may be useful in avoiding excessive nal voltage swings for operation over a large dynamic range.

Consequently, the standard implementations of the photodetector circuit use relatively short gate lengths whose nominal values are bracketed in each wafer fabrication run.

Exact analysis of short channel effects requires numerical simulation since the three dimensional geometry of the device and dopant impurity profiles must be considered. However, a two dimensional model is of some use provided its limitations are understood.

A particularly simple relationship between V_{DS} and the source-drain current of a short channel MOSFET is given in [4],[6]

$$I_S = I_{SL}\exp(qV_{DS}/\delta kT) \quad , \quad (8)$$

where I_{SL} is the source current given by Equation (3) for the long channel MOSFET model and δ is a parameter having a weak dependence on V_{DS} and determined by numerical or empirical means. Note that $\delta = \infty$ corresponds to the long channel model.

Consequently, Equation (7), the explicit temperature dependent for the subthreshold current I_S in the photodetector circuit, can be modified to include the short channel effects of the drain and source voltages

$$I_S(T, V_{BS}, \delta) =$$

$$\frac{\left[GT^{7/2}\exp\left[-E_g/kT+qV_{GS}/kT+q(V_{DB}-V_{BS})/\delta kT-a(T_o)\,(T_o/T)\,(q(V_{GS}+V_{BS})/kT_o)^{\frac{1}{2}}\right]\right]}{\sqrt{2}\,(qV_{GS}/kT+qV_{BS}/kT-a(T_o)\,(T_o/T)\,(q(V_{GS}+V_{BS})/kT_o)^{\frac{1}{2}})^{\frac{1}{2}}}\,, \tag{9}$$

where V_{DB} is the MOSFET drain to substrate voltage.

2.2 Photodiode Model

Under reverse bias conditions, the photodiode current I_D consists of a photocurrent I_p and a dark current I_L. I_p is normally proportional to light intensity and is given by

$$I_p = P\lambda nq/hc\,, \tag{10}$$

where

- n = quantum efficiency
- P = incident illumination power
- λ = wavelength of the light
- h = Planck's constant
- c = velocity of light in free space

Previously reported measurements confirm that the photocurrent has a low sensitivity to temperature and bias voltage [7]-[8].

The photodiode dark current I_L is largely the result of the thermal generation of carriers in the semiconductor. Two components, a generation current produced by the generation of carriers in the reverse biased depletion layer, and a diffusion current caused by the diffusion of carriers generated in the semiconductor bulk to the depletion layer, can be identified. They are given by [8]:

$$I_{gen} = 1/2qn_iW(V)\,A/\tau_o\,, \tag{11}$$

where

- W(V) = depletion layer width proportional to V^x, where V is the reverse bias voltage and $1/3 \le x \le 1/2$
- τ_o = average carrier lifetime

and

$$I_{diff} = qD_n An_i^2/N_A L_n + qD_p An_i^2/N_D L_p \ , \qquad (12)$$

where

D_n = diffusion constant for electrons
D_p = diffusion constant for holes
N_A = dopant impurity concentration for p-type semiconductor material
N_D = dopant impurity concentration for n-type semiconductor material
L_n = diffusion length of electron in p-type semiconductor material
L_p = diffusion length of holes in n-type semiconductor material.

For a photodiode fabricated as a n+ diffusion on a p-type substrate, the second term in Equation (12) can be neglected.

The temperature dependence of Equations (11) and (12) is largely determined by that of the n_i and n_i^2 parameters respectively. Consequently, the temperature dependence of the dark current has the form

$$I_L(T) = I_{gen}(T) + I_{diff}(T) \qquad (13)$$
$$= H_1 \exp(-E_g/2kT) + H_2 \exp(-E_g/kT)$$

For silicon diodes reverse biased by more than a few volts, the I_{gen} component is usually dominant at temperatures up to at least 400°K with the exact crossover temperature being a function of bias voltage and the dopant impurity profile [9]-[10].

Additional mechanisms contributing to dark current such as surface states exist, but these can usually be neglected for high quality silicon photodiodes.

3.0 ANALYSIS OF CIRCUIT MODEL FOR TEMPERATURE DEPENDENCE

Since the steady state photodiode and MOSFET currents are equal, it is possible to set up the equality

$$(I_p + I_L(T)) = I_s(V_{BS}, T, \delta) \ . \qquad (14)$$

This is a useful relationship for analyzing the circuit operation. Since the circuit is operated so that V_{BS} has a substantial DC bias or offset relative to both the drain voltage of the MOSFET and the substrate potential, the temperature dependence of V_{BS} for constant illumination is of critical importance. It determines the need for temperature calibration/compensation techniques and affects the design of the interface electronics which must be able to operate properly over the range of V_{BS} that can result from the worst case combination of illumination and temperature extremes.

The dependence of V_{BS} with respect to temperature is found by evaluating the differential

$$\frac{dV_{BS}}{dT} = \frac{\partial V_{BS}}{\partial (I_s(T, V_{BS}, \delta))} \left(\frac{\partial I_L(T)}{\partial T} - \frac{\partial I_S(T, V_{BS}, \delta)}{\partial T} \right) . \qquad (15)$$

By neglecting the weak temperature dependence of G, Equation (15) can be evaluated to yield

$$\frac{dV_{BS}}{dT} =$$

$$\left(\left(\frac{1}{(I_p + I_L(T))} \frac{\partial I_L}{\partial T} \right) - \left(\frac{7/2}{T} + \frac{1}{T^2} (E_g/k - qV_{GS}/k - q(V_{DB} - V_{BS})/\delta k + a(T_o) T_o (q(V_{GS} + V_{BS})/kT_o)^{\frac{1}{2}} \right) \right.$$

$$\left. \times \frac{-(kT/q) (q(V_{GS} + V_{BS})/kT_o)^{\frac{1}{2}}}{(1/\delta) (q(V_{GS} + V_{BS})/kT_o)^{\frac{1}{2}} + \frac{1}{2} a(T_o)} \right) , \qquad (16)$$

where

$$\frac{\partial I_L}{\partial T} \frac{kT}{qI_S} = Eg/2qT : (I_L \gg I_p) \qquad (17)$$

$$\frac{\partial I_L}{\partial T} \frac{kT}{qI_S} = 0 : (I_L \ll I_p) \qquad (18)$$

From the examination of these results, dV_{BS}/dT has only a mild dependence on temperature and is virtually constant over restricted temperature ranges. They also indicate that dV_{BS}/dT decreases with increasing I_p due to the dependence of V_{BS} on I_p, except for very small values of I_p where I_L is important.

Another important performance parameter is the logarithmic slope of the response given by

$$S_R = \log_{10} \frac{\partial V_{BS}}{\partial \ln I_p} \quad . \tag{19}$$

S_R represents the change in output voltage resulting from a change in I_p by one decade. It is desirable that S_R has a low sensitivity to temperature and I_p, as complex calibration and correction schemes would otherwise be necessary in applications where accurate measurements of illumination intensity are desirable.

Using Equation (9) yields

$$S_R \approx -2.3 \frac{I_p}{(I_p + I_L(T))} \left(\frac{kT}{q}\right) \left(\frac{(q(V_{GS}+V_{BS})/kT_o)^{\frac{1}{2}}}{(1/\delta)(q(V_{GS}+V_{BS})/kT_o)^{\frac{1}{2}} + \frac{1}{2}a(T_o)} \right) \quad . \tag{20}$$

Except for very low values of I_p where the ratio $I_L(T)/I_p$ becomes significant, this result has a weak inverse dependence on I_p since V_{BS} is a logarithmic function of I_p. The temperature dependence of S_R consists of an explicit component linearly proportional to T and implicit non-linear components resulting from the temperature dependence of $I_L(T)$ and V_{BS}.

4.0 EXPERIMENTAL RESULTS

The output voltage from a photodetector structure on the DALSA D4-100 evaluation device was measured for different combinations of illumination intensity and temperature. Illumination was obtained using a HeNe laser having a wavelength of 632.8 nanometers. Uniform illumination over the photodiode and non-critical requirements for its alignment were ensured by the use of a beam expander. An optical attenuator allowed illumination intensity to be varied over a range of 45 dB. The voltage was directly measured by switching on the reset MOSFET and using a low input bias current amplifier ($i_{bias} \approx 10^{-12}$ A) to

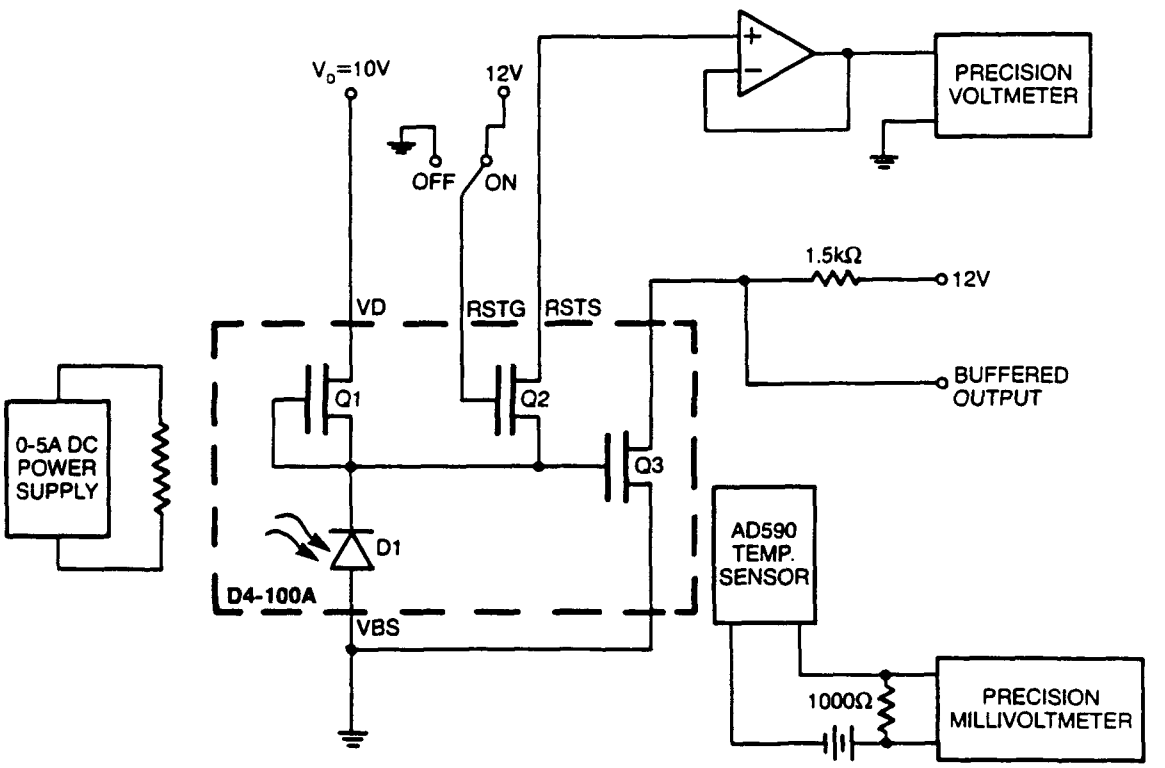

FIGURE 2: EXPERIMENT SET-UP

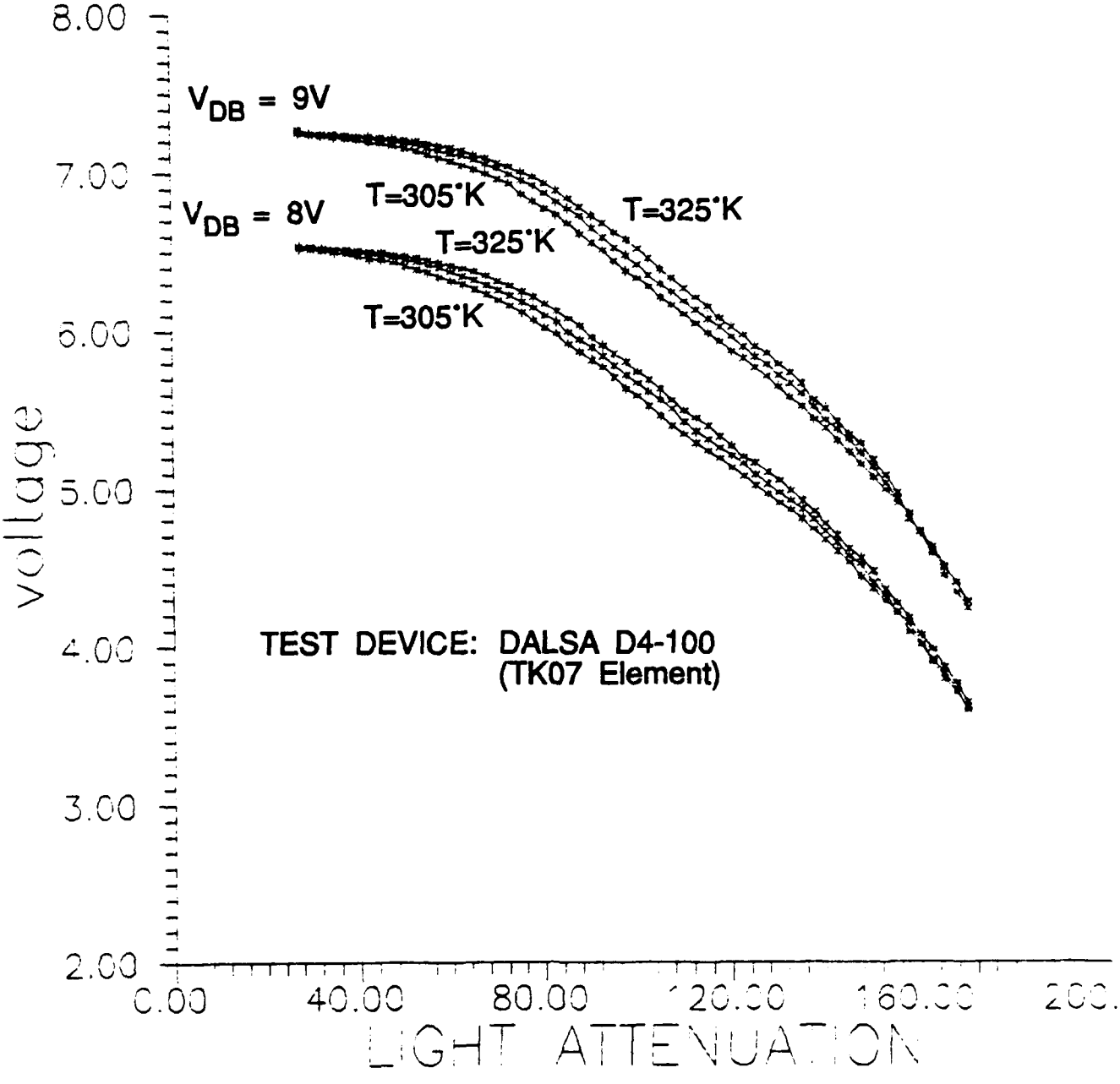

FIGURE 3: V_{BS} (VOLTS) VS LIGHT ILLUMINATION ATTENUATION (INCREMENTS OF 0.3 dB/UNIT)

monitor the voltage at the open circuited reset bias line as shown in Figure 2. Typical results, given in Figure 3, are generally consistent with the analysis.

At 305°K, S_R and dV_{BS}/dT are approximately 0.76V/decade and 10 mV/°K over a large range of illumination levels. Using the a = 1.74 determined from the nominal process parameters [13], Equations (20) and (16) can be fitted to the data corresponding to settings of the optical attenuator between 80 and 140 by setting V_{GS} = 1.6 V and δ = 36 (V_{BS} = 6V).

At low illumination levels S_R and dV_{BS}/dT decrease substantially, as expected from the analysis, but the accuracy of the results is uncertain due to the likelihood that leakage current in the experimental setup is significant.

The dependence of V_{BS} on V_{DB} also demonstrates a good fit to the analysis. The 1 volt change in V_{DB} from 8.0 to 9.0 volts shifts V_{BS} by approximately 0.7 volts for a wide range of illumination levels. This is reasonable since the substitution of the two sets of values for V_{DB} and V_{BS} in Equation (9) results in little change in current.

One area of discrepancy concerns the temperature dependence of S_R which varies by approximately half the amount predicted by the kT/q term in Equation (20) for a 20°C temperature change. However, this may be a result of the limitations of the experimental test setup, which did not permit a large temperature range to be used or an accurate direct measurement of the chip temperature of the photodetector device.

5.0 DISCUSSION

It is interesting to compare the analytical behaviour of the large dynamic range photodetector circuit with that of a photodiode operated in the solar cell mode. The logarithmic slope of the response of the latter can be derived from Equation (1) as,

$$S_R = 2.3 \frac{mkT}{q}\left(\frac{I_p}{I_p+I_o}\right) . \qquad (21)$$

This is similar in form to Equation (20), but differs in magnitude by approximately an order of magnitude due to the difference between m which is between 1 and 2 and the last term in brackets.

The photodiode output voltage V has a temperature dependence given by

$$\frac{dV}{dT} \approx -\frac{1}{T}\left[(3kT/q+mEg/q)\left(\frac{I_p}{I_p+I_o(T)}\right)-V\right] . \qquad (22)$$

Equations (16) and (22) both have a relatively weak ($\approx 1/T$) dependence on temperature and tend to decrease with increasing illumination. The behaviour for the large dynamic range photodetector is more complex in that it is dependent on fabrication process parameters (which affect V_{GS}, $a(T_o)$ and δ), operating conditions (V_{DB}) and the MOSFET geometry (L, dopant profile). If the results of Equations (16) and (22) are normalized with respect to the appropriate S_R, a useful figure of merit is obtained.

$$F = \left[\frac{dV}{dT}\frac{1}{S_R}\right]^{-1} . \qquad (23)$$

F can be regarded as the temperature change corresponding to an output voltage change that would indicate a change in current by an order of magnitude. For the large dynamic range photodetector, this can be evaluated to yield

$$F \approx \frac{2.3I_p}{I_p+I_L(T)}$$

$$\left[\left(7/2_{/T}+\frac{1}{T^2}(E_g/k-qV_{GS}/k-q(V_{DB}-V_{BS})/\delta k-a(T_o)\,T_o(q(V_{GS}+V_{BS})/kT_o)^{\frac{1}{2}})\right)\right] \qquad (24)$$

$$\left[-\frac{1}{(I_p+I_L(T))}\frac{\partial I_L}{\partial T}\right]^{-1} .$$

Similarly, for the photodiode operated in the solar cell mode

$$F \approx \frac{2.3I_p}{I_p+I_o(T)}\frac{mkT^2}{q}\left[(3kT/q+mEg/q)\frac{I_p}{I_p+I_o(T)}-V\right]^{-1} . \qquad (25)$$

For operation at 300°K with moderate illumination, typical values of F would be approximately 68°K and 33°K for Equations (24) and (25) respectively.

A more complete treatment of the problem would involve the generation of a model for δ which would include its voltage dependence, the collection of data from a larger number of sample devices whose process parameters have been independently evaluated and a more accurate test setup as proposed in Appendix 1.

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APPENDIX 1

IMPROVED TEST CIRCUIT FOR DC CHARACTERIZATION OF THE DYNAMIC RANGE PHOTODETECTOR CIRCUIT

The test circuit used in the work described in this report has several deficiencies which result from the constraint of having to use a standard production photodetector device. Usually the photodetector circuit is interfaced to an amplifier which is typically a MOSFET operated as either a common source or common drain amplifier or is multiplexed via a CCD Parallel-Input-Serial-Output shift register. However, it becomes difficult to determine what the actual voltage V_{BS} in the photodetector circuit is since the amplifier transfer function will not be accurately known. Consequently, the direct measurement of the DC voltage from the photodetector must be done via the reset bias voltage input with the reset transistor tuned on. This has the fundamental disadvantage that loading effects will be introduced at low illumination levels. Even if a low input bias current amplifier is used to buffer the signal, as was done for the experimental work described in the report, significant leakage currents can result from contamination of the interconnection between the photodetector and amplifier.

An attractive solution is shown in Figure A-1. This involves the use of a special test device with a pair of matching common drain amplifiers A1 and A2. Amplifier A1 buffers the photodetector voltage and amplifier A2 is used inside the feedback loop of an operational amplifier A3 to compensate for the input to output offset voltage and non-ideal transfer function of amplifier A1. Since the input of amplifier A1 needs not be accessible externally, it has no need for electrostatic protection circuitry which will introduce leakage currents. Consequently, the photodetector circuit operates with negligible DC loading and high accuracy should be achievable.

The test device could also have additional test structures to allow the accurate measurements of the fabrication process parameters.

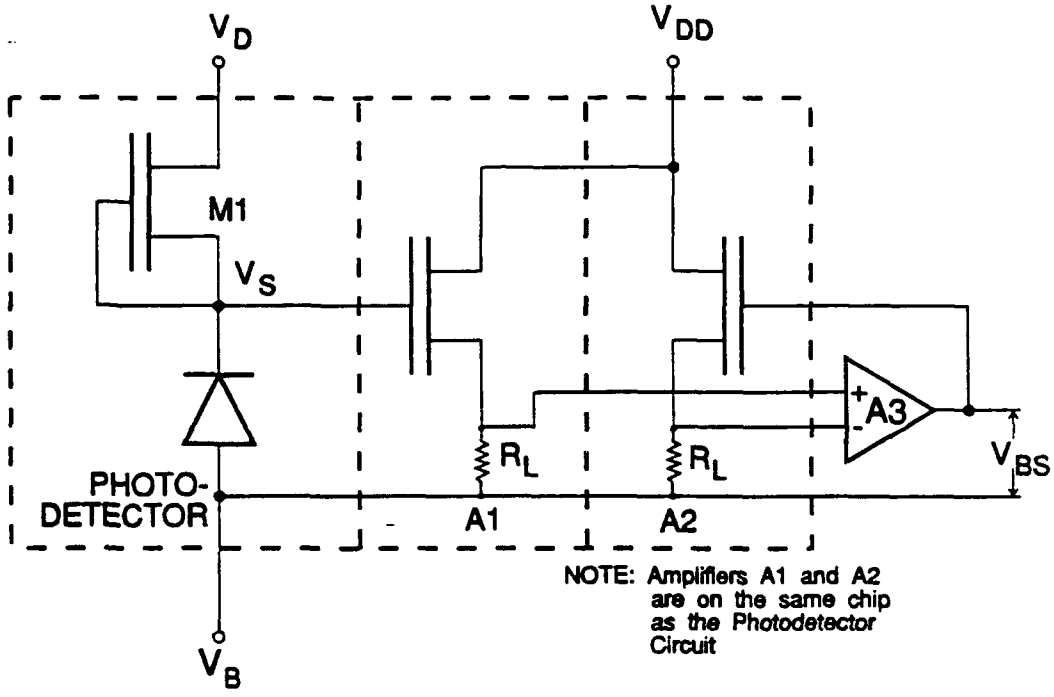

FIGURE A1: IMPROVED PHOTODETECTOR TEST CIRCUIT FOR DC CHARACTERIZATION

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